A PROBABILISTIC PROOF OF WALLIS'S FORMULA FOR π

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There are many beautiful formulas for π (see for example [4]). The purpose of this note is to introduce an alternate derivation of Wallis's product formula, equation (1), which could be covered in a first course on probability, statistics, or number theory. We quickly review other famous formulas for π , recall some needed facts from probability, and then derive Wallis's formula. We conclude by combining some of the other famous formulas with Wallis's formula to derive an interesting expression for $\log(\pi/2)$ (equation (5)).

Often in a first-year calculus course students encounter the Gregory-Leibniz formula,

$$\frac{\pi}{4} = 1 - \frac{1}{3} + \frac{1}{5} - \frac{1}{7} + \dots = \sum_{n=0}^{\infty} \frac{(-1)^n}{2n+1}.$$

The proof uses the fact that the derivative of $\arctan x$ is $1/(1+x^2)$, so $\pi/4=\int_0^1 dx/(1+x^2)$. To complete the proof, expand the integrand with the geometric series formula and then justify interchanging the order of integration and summation.

Another interesting formula involves Bernoulli numbers and the Riemann zeta function. The Bernoulli numbers B_k are the coefficients in the Taylor series

$$\frac{t}{e^t - 1} = 1 - \frac{t}{2} + \sum_{k=2}^{\infty} \frac{B_k t^k}{k!};$$

each B_k is rational. The Riemann zeta function is $\zeta(s) = \sum_{n=1}^{\infty} n^{-s}$, which converges for real part of s greater than 1. Using complex analysis one finds (see for instance [10, p. 365] or [18, pp. 179–180]) that

$$\zeta(2k) = -\frac{(-4)^k B_{2k}}{2 \cdot 2^{k!}} \cdot \pi^{2k},$$

yielding formulas for π to any even power. In particular, $\pi^2/6=\sum_n n^{-2}$ and $\pi^4/90=\sum_n n^{-4}$.

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 $^{^1}$ An amusing consequence of these formulas is a proof of the infinitude of primes. Using unique factorization, one can show that $\zeta(s)$ also equals $\prod_p (1-p^{-s})^{-1}$, where p runs over all primes. As π^2 is irrational and $\zeta(2)=\pi^2/6$, there must be infinitely many primes: if there were only finitely many then $\pi^2/6=\prod_p (1-p^{-2})^{-1}$ would be rational! See [13] for explicit lower bounds on $\pi(x)$ derivable from upper bounds for the irrationality measure of $\zeta(2)$, and [14] for more details on the numerous connections between $\zeta(s)$ and number theory.

One of the most interesting formulas for π is a multiplicative one due to Wallis (1665):

$$\frac{\pi}{2} = \frac{2 \cdot 2}{1 \cdot 3} \frac{4 \cdot 4}{3 \cdot 5} \frac{6 \cdot 6}{5 \cdot 7} \frac{8 \cdot 8}{7 \cdot 9} \cdots = \prod_{n=1}^{\infty} \frac{2n \cdot 2n}{(2n-1)(2n+1)}.$$
 (1)

Common proofs use the infinite product expansion for $\sin x$ (see [18, p. 142]) or induction to prove formulas for integrals of powers of $\sin x$ (see [3, p. 115]). We present a mostly elementary proof using standard facts about probability distributions encountered in a first course on probability or statistics (and hence the title). The reason we must write "mostly elementary" is that at one point we appeal to the Dominated Convergence Theorem. It is possible to bypass this and argue directly, and we sketch the main ideas for the interested reader.

Recall that a continuous function f(x) is a continuous probability distribution if (1) $f(x) \geq 0$ and (2) $\int_{-\infty}^{\infty} f(x) dx = 1$. We immediately see that if g(x) is a nonnegative continuous function whose integral is finite then there exists an a>0 such that ag(x) is a continuous probability distribution (take $a=1/\int_{-\infty}^{\infty} g(x) dx$). This simple observation is a key ingredient in our proof, and is an extremely important technique in mathematics; the proof of Wallis's formula is but one of many applications. In fact, this observation greatly simplifies numerous calculations in random matrix theory, which has successfully modeled diverse systems ranging from energy levels of heavy nuclei to the prime numbers; see [5, 14] for introductions to random matrix theory and [11] for applications of this technique to the subject. One of the purposes of this paper is to introduce students to the consequences of this simple observation.

Our proof relies on two standard functions from probability, the Gamma function and the Student t-distribution. The Gamma function $\Gamma(x)$ is defined by

$$\Gamma(x) = \int_0^\infty e^{-t} t^{x-1} dt.$$

Note that this integral is well defined if the real part of x is positive. Integrating by parts yields $\Gamma(x+1)=x\Gamma(x)$. This implies that if n is a nonnegative integer then $\Gamma(n+1)=n!$; thus the Gamma function generalizes the factorial function (see [17] for more on the Gamma function, including another proof of Wallis's formula involving the Gamma function). We need the following:

Claim:
$$\Gamma(1/2) = \sqrt{\pi}$$
.

 $^{^2}$ For a statistical proof involving an experiment and data, see the chapter on Buffon's needle in [1] (page 133): if you have infinitely many parallel lines d units apart, then the probability that a "randomly" dropped rod of length $\ell \leq d$ crosses one of the lines is $2\ell/\pi d$. Thus you can calculate π by throwing many rods on the grid and counting the number of intersections.

 $^{^{3}}$ A nice application of Wallis's formula is in determining the universal constant in Stirling's formula for n!; see [15] for some history and applications.

Proof. In the integral for $\Gamma(1/2)$, change variables by setting $u = \sqrt{t}$ (so $dt = 2udu = 2\sqrt{t}du$). This yields

$$\Gamma(1/2) = 2 \int_0^\infty e^{-u^2} du = \int_{-\infty}^\infty e^{-u^2} du.$$

This integral is well-known to equal $\sqrt{\pi}$ (see page 542 of [2]). The standard proof is to square the integral and convert to polar coordinates:

$$\Gamma(1/2)^2 = \int_{-\infty}^{\infty} e^{-u^2} du \int_{-\infty}^{\infty} e^{-v^2} dv = \int_{0}^{\infty} \int_{0}^{2\pi} e^{-r^2} r dr d\theta = \pi.$$

In fact, our proof above shows

$$\int_{-\infty}^{\infty} \frac{1}{\sqrt{2\pi}} e^{-t^2/2} dt = 1.$$
 (2)

This density is called the standard normal (or Gaussian). This is one of the most important probability distributions, and we shall see it again when we look at the Student t-distribution. If g is a continuous probability density, then we say that the random variable Y has distribution g if for any interval [a,b] the probability that Y takes on a value in [a,b] is $\int_a^b g(y)dy$. The celebrated Central Limit Theorem (see [6, p. 515] for a proof) states that for many continuous densities g, if X_1, \ldots, X_n are independent random variables, each with density g, then as $n \to \infty$ the distribution of $(Y_n - \mu)/\sigma$ converges to the standard normal (where $Y_n = (X_1 + \cdots + X_n)/n$ is the sample average, μ is the mean of g, and σ is its standard deviation⁴).

The second function we need is the Student⁵ t-distribution (with ν degrees of freedom):

$$f_{\nu}(t) = \frac{\Gamma\left(\frac{\nu+1}{2}\right)}{\sqrt{\pi\nu} \Gamma\left(\frac{\nu}{2}\right)} \cdot \left(1 + \frac{t^2}{\nu}\right)^{-\frac{\nu+1}{2}} = c_{\nu} \left(1 + \frac{t^2}{\nu}\right)^{-\frac{\nu+1}{2}};$$

here ν is a positive integer and t is any real number.

Claim: The Student t-distribution is a continuous probability density.

Proof. As $f_{\nu}(t)$ is clearly continuous and nonnegative, to show $f_{\nu}(t)$ is a probability density it suffices to show that it integrates to 1. We must therefore show that

$$\int_{-\infty}^{\infty} \left(1 + \frac{t^2}{\nu} \right)^{-\frac{\nu+1}{2}} dt = \frac{\sqrt{\pi\nu} \Gamma\left(\frac{\nu}{2}\right)}{\Gamma\left(\frac{\nu+1}{2}\right)}.$$

⁴The mean μ of a distribution is its average value: $\mu = \int xg(x)dx$. The standard deviation σ measures how spread out a distribution is about its average value: $\sigma^2 = \int (x - \mu)^2 g(x) dx$.

⁵Student was the pen name of William Gosset.

As the integrand is symmetric, we may integrate from 0 to infinity and double the result. Letting $t = \sqrt{\nu} \tan \theta$ (so $dt = \sqrt{\nu} \sec^2 \theta d\theta$) we find

$$\int_{-\infty}^{\infty} \left(1 + \frac{t^2}{\nu} \right)^{-\frac{\nu+1}{2}} dt = 2\sqrt{\nu} \int_{0}^{\pi/2} \frac{\sec^2 \theta d\theta}{\sec^{\nu+1} \theta} = 2\sqrt{\nu} \int_{0}^{\pi/2} \cos^{\nu-1} \theta d\theta.$$

The proof follows immediately from two properties of the Beta function (see [2, p. 560]):

$$B(p,q) = \Gamma(p)\Gamma(q)/\Gamma(p+q)$$

$$B(m+1,n+1) = 2\int_0^{\pi/2} \cos^{2m+1}(\theta)\sin^{2n+1}(\theta)d\theta;$$
(3)

an elementary proof without appealing to properties of the Beta function is given in Appendix A. \Box

The Student t-distribution arises in statistical analyses where the sample size ν is small and each observation is normally distributed with the same mean and the same (unknown) variance (see [8, 12]). The reason the Student t-distribution is used only for small samples sizes is that as $\nu \to \infty$, $f_{\nu}(t)$ converges to the standard normal; proving this will yield Wallis's formula. While we can prove this by invoking the Central Limit Theorem, we may also see this directly by recalling that

$$e^x = \lim_{N \to \infty} \left(1 + \frac{x}{N} \right)^N.$$

We therefore have

$$\lim_{\nu \to \infty} \left(1 + \frac{t^2}{\nu} \right)^{-\nu/2} = \left(e^{t^2} \right)^{-1/2} = e^{-t^2/2}.$$

As $f_{\nu}(t)$ is a probability distribution for all positive integers ν , it integrates to 1 for all such ν , which is equivalent to

$$\frac{1}{c_{\nu}} = \frac{\sqrt{\pi\nu} \Gamma\left(\frac{\nu}{2}\right)}{\Gamma\left(\frac{\nu+1}{2}\right)} = \int_{-\infty}^{\infty} \left(1 + \frac{t^2}{\nu}\right)^{-\frac{\nu+1}{2}} dt.$$

Taking the limit as $\nu \to \infty$ yields

$$\lim_{\nu \to \infty} \frac{1}{c_{\nu}} = \lim_{\nu \to \infty} \int_{-\infty}^{\infty} \left(1 + \frac{t^{2}}{\nu} \right)^{-\frac{\nu+1}{2}} dt$$

$$= \int_{-\infty}^{\infty} \lim_{\nu \to \infty} \left(1 + \frac{t^{2}}{\nu} \right)^{-\frac{\nu+1}{2}} dt = \int_{-\infty}^{\infty} e^{-t^{2}/2} dt = \sqrt{2\pi}.$$

Some work is necessary of course to justify interchanging the integral and the limit; this justification is why our argument is only "mostly elementary". A standard proof uses the Dominated Convergence Theorem (see [7, p. 54] or [9, p. 238]) to

show that as $\nu\to\infty$ the t-distribution converges to the standard normal; one may take $2008\exp(-t^2/2008)$ as the dominating function. We have therefore shown that

$$c = \lim_{\nu \to \infty} c_{\nu} = \lim_{\nu \to \infty} \frac{\Gamma\left(\frac{\nu+1}{2}\right)}{\sqrt{\pi\nu} \Gamma\left(\frac{\nu}{2}\right)} = \frac{1}{\sqrt{2\pi}}.$$
 (4)

The fact that $c=1/\sqrt{2\pi}$ is the key step in our proof of Wallis's formula. We have calculated the limit by using our observation that a probability distribution must integrate to 1; calculating it by brute force analysis of the Gamma factors yields our main result.

Theorem: Wallis's formula is true.

Proof. The proof follows from expanding the Gamma functions and substituting into (4); we highlight the main steps. Let $\nu=2m$. Using $\Gamma(n+1)=n\Gamma(n)$ and $\Gamma(1/2)=\sqrt{\pi}$ we find that

$$\Gamma\left(\frac{2m+1}{2}\right) = (2m-1)(2m-3)\cdots 5\cdot 3\cdot 1\cdot \sqrt{\pi}/2^m.$$

As $\Gamma(m) = (m-1)!$, after some algebra we find that

$$c_{2m} = \frac{1 \cdot 3 \cdot 5 \cdots (2m-3) \cdot (2m-1)}{2 \cdot 4 \cdot 6 \cdots (2m-2) \cdot 2m} \frac{\sqrt{m}}{\sqrt{2}}.$$

Multiplying by $1 \cdot (2m+1)/(2m+1)$ and regrouping, we find

$$c_{2m} = \left(\frac{1\cdot 3}{2\cdot 2} \cdot \frac{3\cdot 5}{4\cdot 4} \cdots \frac{(2m-1)(2m+1)}{2m\cdot 2m} \frac{1}{2m+1}\right)^{\frac{1}{2}} \frac{\sqrt{m}}{\sqrt{2}},$$

which we rewrite as

$$\prod_{n=1}^{m} \frac{2n \cdot 2n}{(2n-1)(2n+1)} = \frac{2 \cdot 2}{1 \cdot 3} \frac{4 \cdot 4}{3 \cdot 5} \cdots \frac{2m \cdot 2m}{(2m-1)(2m+1)} = \frac{m}{(4m+2)c_{2m}^2}.$$

As $\lim_{m\to\infty} 1/c_{2m}^2 = 2\pi$ and $\lim_{m\to\infty} m/(4m+2) = 1/4$, we have

$$\prod_{n=1}^{\infty} \frac{2n \cdot 2n}{(2n-1)(2n+1)} \; = \; \frac{\pi}{2},$$

which completes the proof.

By combining the expansion for π from Wallis's formula with those involving the Bernoulli numbers and the zeta function, we obtain a proof of the following amusing formula for $\log(\pi/2)$.

⁶For completeness we sketch how such an argument could proceed. If $t \in [-\log^2 \nu, \log^2 \nu]$ then $|(1+t^2/\nu)^{-(\nu+1)/2} - \exp(-t^2/2)|$ tends to zero rapidly with ν . Further, if f(t) is the density of the standard normal, then $\int_{|t| \ge \log^2 \nu} f(t) dt$ and $\int_{|t| \ge \log^2 \nu} f_{\nu}(t) dt$ also tend to zero rapidly with ν . Careful bookkeeping shows that the normalization constants c_{ν} must therefore approach $c = 1/\sqrt{2\pi}$.

Theorem: We have

$$\log \frac{\pi}{2} = \sum_{k=1}^{\infty} \frac{\zeta(2k)}{4^k \cdot k} = -\sum_{k=1}^{\infty} \frac{(-1)^k B_{2k}}{2k \cdot 2k!} \cdot \pi^{2k}.$$
 (5)

Proof. The Taylor series of $\log(1-x)$ is $-\sum_{k=1}^{\infty}x^k/k$. The nth factor in Wallis's formula may be written as $\left(1-\frac{1}{4n^2}\right)^{-1}$. Thus taking logarithms of Wallis's formula and Taylor expanding yields

$$\log \frac{\pi}{2} = -\sum_{n=1}^{\infty} \log \left(1 - \frac{1}{4n^2} \right) = \sum_{n=1}^{\infty} \sum_{k=1}^{\infty} \frac{1}{k \cdot (4n^2)^k} = \sum_{k=1}^{\infty} \frac{1}{4^k \cdot k} \sum_{n=1}^{\infty} \frac{1}{n^{2k}}.$$

The *n*-sum gives $\zeta(2k)$, and the claim now follows.

Note that the above formula for $\log(\pi/2)$ converges well. It is easy to see that $|\zeta(2k)| \leq 2$ (and $\lim_{k \to \infty} \zeta(2k) = 1$). Thus each additional summand yields at least one new digit (base 4). See [16] for additional formulas for $\log(\pi/2)$.

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APPENDIX A. ELEMENTARY CALCULATION OF CONSTANTS

For completeness, we provide a more elementary derivation that the stated constant is the correct normalization constant for the Student t-distribution.

Lemma A.1. We have

$$2\sqrt{\nu} \int_0^{\pi/2} \cos^{\nu-1} \theta d\theta = \frac{\sqrt{\pi\nu} \Gamma\left(\frac{\nu}{2}\right)}{\Gamma\left(\frac{\nu+1}{2}\right)}.$$
 (6)

Proof. The claim follows by induction; we sketch the main idea. Assume we have proven the claim for all $\nu \leq n$. Then

$$\int_0^{\pi/2} \cos^{n+1} \theta d\theta = \int_0^{\pi/2} (1 - \sin^2 \theta) \cos^{n-1} \theta d\theta$$
$$= \int_0^{\pi/2} \cos^{n-1} \theta d\theta - \int_0^{\pi/2} \sin \theta \cdot (\cos^{n-1} \theta \sin \theta) d\theta.$$

We integrate the second term on the right by parts, with $u = \sin \theta$ and $dv = \cos^{n-1} \theta \sin \theta d\theta$. The uv term vanishes at 0 and $\pi/2$ and we are left with

$$\int_0^{\pi/2} \cos^{n+1} \theta d\theta = \int_0^{\pi/2} \cos^{n-1} \theta d\theta - \frac{1}{n} \int_0^{\pi/2} \cos^n \theta \cos \theta d\theta,$$

which simplifies to

$$\int_0^{\pi/2} \cos^{n+1} \theta d\theta = \frac{n}{n+1} \int_0^{\pi/2} \cos^{n-1} \theta d\theta.$$

The claim now follows from standard properties of the Gamma function ($\Gamma(m+1)=m\Gamma(m)$ and $\Gamma(1/2)=\sqrt{\pi}$).

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