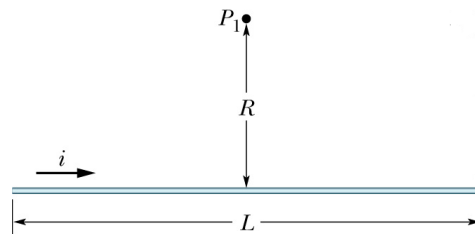
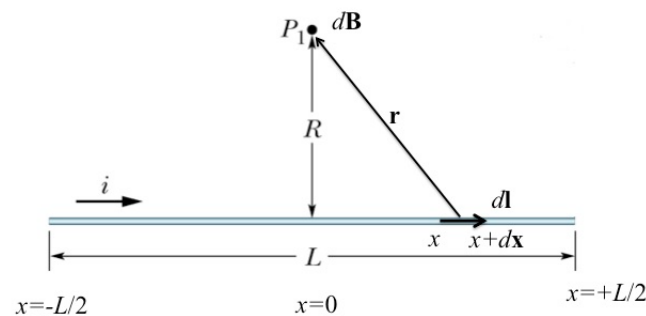


1. Show, that at point P_1 the magnetic field is given by: $B = \frac{\mu_0 i}{2\pi R} \frac{L}{(L^2 + 4R^2)^{1/2}}$. What happens if $L \rightarrow \infty$?



Solution:



To solve this problem we need to apply Biot-Savart Law. We consider the elementary part $d\mathbf{l}$ of the wire at a position x having length dx . Thus $d\mathbf{l} = dx\hat{\mathbf{x}}$. This part is flown by a current I so at the point P_1 it creates a magnetic field $d\mathbf{B}$ given by:

$$d\mathbf{B} = \frac{\mu_0 I}{4\pi} \cdot \frac{d\mathbf{l} \times \mathbf{r}}{r^3}$$

where \mathbf{r} is a vector having its tail (beginning) at the tail of $d\mathbf{l}$ and its tip (end) at the point P_1 . Thus $\mathbf{r} = (0, h, 0) - (x, 0, 0)$ or $\mathbf{r} = (-x, h, 0)$. Then

$$d\mathbf{l} \times \mathbf{r} = \begin{vmatrix} \hat{\mathbf{x}} & \hat{\mathbf{y}} & \hat{\mathbf{z}} \\ dx & 0 & 0 \\ 0 & h & -x \end{vmatrix} = R dx \hat{\mathbf{z}}$$

The magnitude of r is given by $r = (x^2 + R^2)^{1/2}$. Thus for the elementary magnetic field we have:

$$d\mathbf{B} = \frac{\mu_0 I}{4\pi} \cdot \frac{R dx}{(x^2 + R^2)^{3/2}} \hat{\mathbf{z}}$$

The total magnetic field is taken by an integration we get

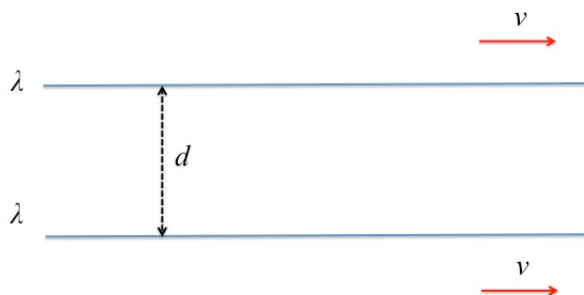
$$\mathbf{B} = \int_{-L/2}^{+L/2} d\mathbf{B} = \hat{\mathbf{z}} \frac{\mu_0 IR}{4\pi} \cdot \int_{-L/2}^{L/2} \frac{dx}{(x^2 + R^2)^{3/2}} = \hat{\mathbf{z}} \frac{\mu_0 IR}{4\pi} \cdot \left[\frac{x}{R^2 (x^2 + R^2)^{1/2}} \right]_{-L/2}^{+L/2}$$

$$\mathbf{B} = \hat{\mathbf{z}} \frac{\mu_0 I}{4\pi} \cdot \frac{L}{R \sqrt{1 + (L/2R)^2}}$$

If L goes to infinity then

$$\mathbf{B} \approx \hat{\mathbf{z}} \frac{\mu_0 I}{4\pi} \cdot \frac{L}{r^2 (\cos \theta)^{3/2}} = \hat{\mathbf{z}} \frac{\mu_0 I}{2\pi R}$$

2. Suppose you have two infinite long straight wires carrying a charge with a linear charge density λ . The two wires move at constant speed v as shown in figure.
- What is the magnetic force per unit length on each wire? (10 marks)
 - What is the electric force on each wire? (6 marks)
 - How great would v have to be in order for the magnetic attraction to balance the electric repulsion? Work out the actual number. Is it a reasonable sort of speed? (4 marks)



You are given that the magnetic field at a distance s from an infinitely current carrying wire is:

$$B = \frac{\mu_0 I}{2\pi s}$$

and the electric field at a distance s from an infinitely charged wire is

$$E = \frac{\lambda}{2\pi\epsilon_0 s}$$

Solution

- i) The moving charged wire is equivalent to a current

$$I_1 = \frac{\Delta q}{\Delta t} = \frac{\lambda \Delta x}{\Delta t} = \lambda v$$

Thus each wire generates a magnetic field at a distance d away from it:

$$B_1 = \frac{\mu_0 \lambda v}{2\pi d}$$

This field is responsible for generating a force on the other wire given by:

$$\Delta F_{2,m} = B_1 I_2 \Delta x = \frac{\mu_0 \lambda v}{2\pi d} I_2 \Delta x = \frac{\mu_0 \lambda^2 v^2}{2\pi d} \Delta x$$

Thus the magnetic force per unit length is:

$$F_m = \frac{\Delta F_{2,m}}{\Delta x} = \frac{\mu_0 \lambda^2 v^2}{2\pi d}$$

- ii) The electric force per unit length on the second wire due to the electric field of the first is:

$$F_{el} = \frac{\Delta F_{2,el}}{\Delta x} = \frac{E \Delta q}{\Delta x} = \frac{\lambda}{2\pi \epsilon_0 d \Delta x} \lambda \Delta x = \frac{\lambda^2}{2\pi \epsilon_0 d}$$

- iii) The magnetic attraction per unit length is given by:

$$F_m = \frac{\mu_0 \lambda^2 v^2}{2\pi d}$$

The electric repulsion per unit length on one wire is:

$$F_{el} = \frac{1}{2\pi \epsilon_0} \frac{\lambda^2}{d}$$

The two forces balance when:

$$F_m = F_{el}$$

$$\frac{\mu_0 \lambda^2 v^2}{2\pi d} = \frac{1}{2\pi \epsilon_0} \frac{\lambda^2}{d}$$

$$v^2 = \frac{1}{\mu_0 \epsilon_0} = c^2$$

$$v = c$$

This is impossible

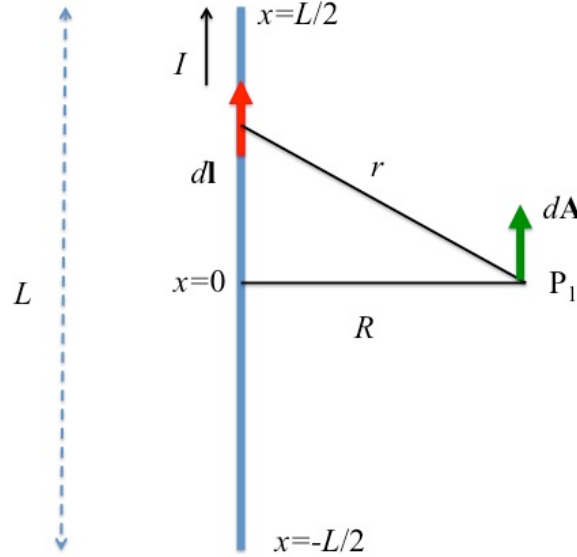
3. Assume you have a thin wire with current I . The vector potential $d\mathbf{A}$ created from an element of length $d\mathbf{l}$ of this wire at a point which is at a distance r from the element is given by the relation:

$$d\mathbf{A} = \frac{\mu_0 I}{4\pi} \frac{d\mathbf{l}}{r}$$

- (i) Use this relation to calculate the vector potential at point P_1 in the figure of problem 1. Assume that the distance R is far smaller than the length of the wire, i.e. $R \ll L$.
- (ii) What will be the expression for the vector potential if the wire has infinite length?

You are given: $(1+x)^{1/2} \approx 1 + \frac{1}{2}x$, for $x \ll 1$, $\int \frac{dx}{(a^2 + x^2)^{1/2}} = \ln \left[x + \sqrt{x^2 + a^2} \right] + C$

Solution:



We consider the elementary part $d\mathbf{l}$ of the wire at a position x having length dx . Thus $d\mathbf{l} = dx\hat{\mathbf{x}}$.

This part is flown by a current I so at the point P_1 it creates a vector potential $d\mathbf{A}$ given by:

$$d\mathbf{A} = \frac{\mu_0 I}{4\pi} \frac{d\mathbf{l}}{r} \Rightarrow \mathbf{A} = \left(\frac{\mu_0 I}{4\pi} \int_{-L/2}^{+L/2} \frac{dx}{r} \right) \hat{\mathbf{x}} \Rightarrow \mathbf{A} = \left(\frac{\mu_0 I}{4\pi} \int_{-L/2}^{+L/2} \frac{dx}{(R^2 + x^2)^{1/2}} \right) \hat{\mathbf{x}}$$

$$\mathbf{A} = \left(\frac{\mu_0 I}{4\pi} \ln \left[x + \sqrt{x^2 + R^2} \right] \right)_{-L/2}^{L/2} \hat{\mathbf{x}} = \frac{\mu_0 I}{4\pi} \left[\ln \left(\frac{L}{2} + \sqrt{\left(\frac{L}{2} \right)^2 + R^2} \right) - \ln \left(-\frac{L}{2} + \sqrt{\left(-\frac{L}{2} \right)^2 + R^2} \right) \right] \hat{\mathbf{x}}$$

$$\mathbf{A} = \frac{\mu_0 I}{4\pi} \left[\ln \left(\frac{L}{2} + \sqrt{\frac{L^2}{4} + R^2} \right) - \ln \left(-\frac{L}{2} + \sqrt{\frac{L^2}{4} + R^2} \right) \right] \hat{\mathbf{x}}$$

Now since $R \ll L$ we have:

$$\mathbf{A} = \frac{\mu_0 I}{4\pi} \left[\ln \left(\frac{L}{2} + \frac{L}{2} \sqrt{1 + \frac{4R^2}{L^2}} \right) - \ln \left(-\frac{L}{2} + \frac{L}{2} \sqrt{1 + \frac{4R^2}{L^2}} \right) \right] \hat{\mathbf{x}} \Rightarrow$$

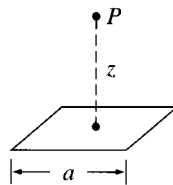
$$\mathbf{A} = \frac{\mu_0 I}{4\pi} \left[\ln \left(\frac{L}{2} + \frac{L}{2} \left(1 + \frac{2R^2}{L^2} \right) \right) - \ln \left(-\frac{L}{2} + \frac{L}{2} \left(1 + \frac{2R^2}{L^2} \right) \right) \right] \hat{\mathbf{x}} \Rightarrow$$

$$\mathbf{A} = \frac{\mu_0 I}{4\pi} \left[\ln \left(L + \frac{R^2}{L} \right) - \ln \left(\frac{R^2}{L} \right) \right] \hat{\mathbf{x}} \Rightarrow \mathbf{A} = \frac{\mu_0 I}{4\pi} \ln \left[\frac{\left(L + \frac{R^2}{L} \right)}{\left(\frac{R^2}{L} \right)} \right] \hat{\mathbf{x}}$$

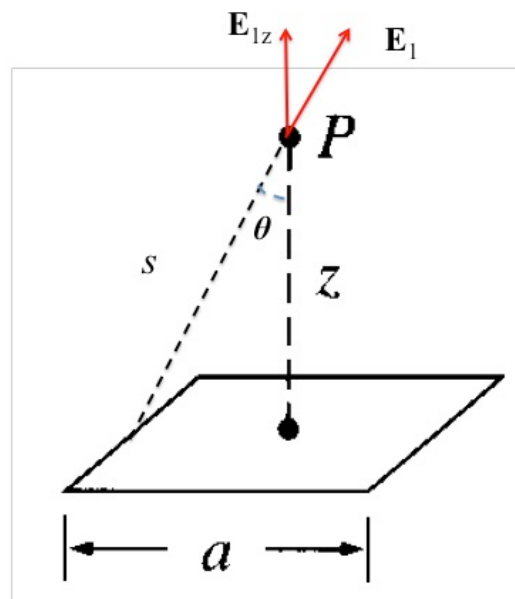
$$\mathbf{A} = \frac{\mu_0 I}{4\pi} \ln \left(\frac{L^2}{R^2} + 1 \right) \Rightarrow_{L \gg R} \mathbf{A} \approx \frac{\mu_0 I}{4\pi} \ln \left(\frac{L^2}{R^2} + 1 \right) \Rightarrow \mathbf{A} \approx \frac{\mu_0 I}{4\pi} 2 \ln \left(\frac{L}{R} \right)$$

$$\mathbf{A} \approx \frac{\mu_0 I}{2\pi} \ln \left(\frac{L}{R} \right)$$

4. Find the electric field (magnitude and direction) a distance z above the center of a square loop as shown in the figure, which carries a uniform line charge λ . (Hint: Use problem 2.3 above).



Solution:



From solution of problem 2.3, the electric field **from one side** on point P is given by:

$$E_1 = \frac{1}{4\pi\epsilon_0} \frac{\lambda a}{s\sqrt{s^2 + a^2/4}} = \frac{1}{4\pi\epsilon_0} \frac{\lambda a}{s\sqrt{s^2 + a^2/4}}$$

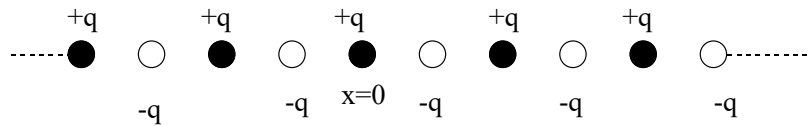
But the total field at P will be made up only from the components along the z -direction so:

$$E_p = 4E_{1z} = 4E_1 \cos \theta = 4 \frac{1}{4\pi\epsilon_0} \frac{\lambda a}{s\sqrt{s^2 + a^2/4}} \frac{z}{s}$$

$$E_p = \frac{1}{\pi\epsilon_0} \frac{\lambda a z}{(z^2 + a^2/4)\sqrt{z^2 + a^2/2}}$$

5. On a straight line we place alternatively an infinite number of charges $+q$ and $-q$ at equal distance as shown in the figure. What is the potential energy of a charge $+q$? You are given that

$$\ln(1+x) = x - \frac{x^2}{2} + \frac{x^3}{3} - \frac{x^4}{4} + \dots + (-1)^{n+1} \frac{x^n}{n} + \dots \quad (-1 < x \leq 1)$$



Solution:

Let the distance between two adjacent charges be d . The potential energy of the charge $+q$ at $x = 0$ is given by

$$U = \frac{1}{4\pi\epsilon_0} \left(-\frac{q^2}{d} - \frac{q^2}{d} + \frac{q^2}{2d} + \frac{q^2}{2d} - \frac{q^2}{3d} - \frac{q^2}{3d} + \dots \right) = \frac{1}{4\pi\epsilon_0} \left(-\frac{2q^2}{d} + \frac{2q^2}{2d} - \frac{2q^2}{3d} + \dots \right)$$

$$= -\frac{2q^2}{4d\pi\epsilon_0} \left(1 - \frac{1}{2} + \frac{1}{3} - \dots \right) = -\frac{q^2}{2d\pi\epsilon_0} \left(1 - \frac{1}{2} + \frac{1}{3} - \dots \right) = -\frac{q^2}{2d\pi\epsilon_0} \ln(1+1) = -\frac{q^2}{2d\pi\epsilon_0} \ln 2$$