PHY 411-506 Computational Physics II

Chapter 8: Statistical Mechanics, Phase Transitions, and the Ising Model

Lecture 1

Monday January 14, 2008

Lecture Outline

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The Ising Model and Statistical Mechanics

Magnetism

• Magnetic Fields are produced by electric currents and magnetic moments

 $\mathbf{B}=\mu_0(\mathbf{H}+\mathbf{M})$

where ${\bf M}$ is the Magnetization and

$$\mathbf{J} =
abla imes \mathbf{H}$$

is the electric current density

• The magnetic susceptibility of the material

$$\chi = \frac{\mathbf{B}}{\mu_0 \mathbf{H}}$$

• Diamagnetism – substance is repelled by a magnetic field

 $\chi < 0$

♦ Atoms have no net magnetic moment

 \diamond Consequence of Lenz' Law – induced moment reduces ${\bf B}$

Magnetism

Paramagnetism – substance is not magnetic and attracted by a magnetic field

 $\chi > 0$

♦ Atoms have net orbital and/or spin magnetic moment

 \diamond Thermal motion causes $\mathbf{M}=\mathbf{0}$

 \diamond Moments tend to line up with ${\bf H}$ – increases ${\bf B}$

- Ferromagnetism substance is magnetic and attracted by a magnetic field
 - \diamond Quantum mechanical Exchange Interactions causes atomic moments to spontaneously align inside domains at temperatures below the Curie temperature $T_{\rm C}$
 - \diamond Magnetic hysteresis relation between ${\bf H}$ and ${\bf B}$ is not linear and depends on history of sample
- Magnetism is an inherently quantum phenomenon

Classical charged particles in thermal equilibrium are not diamagnetic
 Alignment
 Ali

◊ Diamagnetic field of gyrating electrons exactly cancelled by boundary effects – physical argument by Niels Bohr

◇ Proved by H.-J. Van Leeuwen, J. Phys. Radium, 2, 361 (1921).

Ising Model

- Magnetism in metals can be caused by local moments (orbital or spin) and by non-localized moments in electronic bands
- The Heisenberg Model assumes localized moments due to spins s_i at fixed lattice sites i and energy

$$E = -\sum_{i,j} J_{ij} \mathbf{s}_i \cdot \mathbf{s}_j - \mu \sum_i \mathbf{H}_i \cdot \mathbf{s}_i$$

where J_{ij} represents the exchange interaction between spins at sites i, j, μ is the gyromagnetic ratio (magnetic moment) of the spin and \mathbf{H}_i is the external magnetic field at site i

• The Ising Model assumes that the spins are classical and restricted to values

$$s_i = \pm 1$$

Ising Model

and that

$$J_{ij} = \begin{cases} J & \text{for nearest neighbors } \langle ij \rangle \\ 0 & \text{otherwise} \end{cases}$$

so that

$$E = -J\sum_{\langle ij\rangle} s_i s_j - \mu H \sum_i s_i$$

• The Ising Model on a one-dimensional lattice was given by Wilhelm Lenz to Ernst Ising as a PhD thesis topic and has an interesting history

 \diamond lsing found that $T_c = 0$ for this model – it is not ferromagnetic!

• The Ising Model on a 2-d square lattice is ferromagnetic at low temperatures

 \diamond Kramers and Wannier found that

$$\frac{k_{\rm B}T_{\rm c}}{J} = \frac{2}{\log(1+\sqrt{2})} = 2.269\dots$$

 Lars Onsager and C.N. Yang found an exact analytic formula for the magnetization

Statistical Mechanics

- The Ising Model energy (Hamiltonian) has no kinetic term the energy is all potential
- The model has no dynamics cannot study it as a function of time t
- Study its equibrium statistical mechanics
 - ◊ System has a fixed volume spins are at fixed lattice sites
 - \diamond Fix the temperature T by placing it in contact with a large heat reservoir at constant temperature
 - \circ Individual spins can flip by exchanging energy with the reservoir
 - \circ The total energy of the system is not constant
 - \diamond An external magnetic field H is applied to each spin
- \bullet The Canonical Ensemble is a set of configurations or $\it microstates \, \alpha$ of the system with probability

$$P_{\alpha} \sim e^{-E_{\alpha}/(k_{\rm B}T)}$$

 \diamond Observables are averages over this ensemble – for example the

magnetization

$$M = \sum_{\alpha} P_{\alpha} M_{\alpha}$$

where $M_{\alpha} = \sum_{i} s_{i}$

 \bullet For an Ising Model with N spins, there are 2^N microstates or configurations

 \diamond For the 2-d model with $N=20\times20=400$ spins

No. of configs =
$$2^{N_s} = 2^{400} = 2.58 \times 10^{120}$$

To sum all configurations at 1 billion per second would take 2.58×10^{111} seconds = 8.8×10^{103} years

 If analytic solutions are not available, random sampling of configurations (Monte Carlo method) must be used

Mean Field Theory

- Mean Field Theory is an approximate way of finding analytic solutions
- The basic idea is to reduce a many-particle problem to an effective one-particle problem
 - ◊ The spin-spin interaction terms are replaced by an effective magnetic field

$$-J\sum_{\langle ij\rangle}s_is_j \equiv -\mu H_{\rm eff}\sum_i s_i$$

 \diamond Compute the effective field approximately by replacing s_j by the thermal average

$$\langle s_j \rangle \equiv \langle s \rangle$$

 \diamond The

$$H_{\rm eff} = \frac{J}{\mu} \sum_{\rm nearest \ neighbors} \langle s \rangle = \frac{zJ}{\mu} \langle s \rangle$$

where z is the number of nearest neighbors or coordination number of the lattice (z = 2d for a d-dimensional hypercubic lattice)

• The one-spin problem can easily be solved because there are only two configurations $s = \pm 1$ with canonical ensemble probabilities

$$P_{\pm} \sim e^{\pm \mu (H_{\text{eff}} + H)/(k_{\text{B}}T)}$$

 \diamond The thermal average of this spin is

$$\langle s \rangle = \frac{e^{\mu (H_{\rm eff} + H)/(k_{\rm B}T)} - e^{-\mu (H_{\rm eff} + H)/(k_{\rm B}T)}}{e^{\mu (H_{\rm eff} + H)/(k_{\rm B}T)} + e^{-\mu (H_{\rm eff} + H)/(k_{\rm B}T)}} = \tanh\left(\frac{zJ\langle s \rangle + \mu H}{k_{\rm B}T}\right)$$

 \bullet This is an implicit equation for $\langle s \rangle$ – solving for $\langle s \rangle$ is a root finding problem

$$f(\langle s \rangle) \equiv \langle s \rangle - \tanh\left(\frac{zJ\langle s \rangle + \mu H}{k_{\rm B}T}\right) = 0$$

 \bullet For H=0 the equation has a non-zero solution for

$$\frac{k_{\rm B}T}{zJ} \le 1$$

This means the system is ferromagnetic below a Curie temperature

$$T_{\rm c} = \frac{zJ}{k_{\rm B}}$$

- \diamond Wrong for the 1-d model (z = 2)
- ◇ Too high for the 2-d model
- \bullet The spontaneous magnetization $\langle s \rangle$ can be found numerically Exercise 8.1

Root Finding Algorithms

- Finding zeros (roots) of a function needed in many applications see Appendix B of the textbook
- Numerical Recipes has Chapter 9 on Root Finding and Nonlinear Sets of Equations
- Finding extrema of f(x) equivalent to finding roots of f'(x) see Chapter 10 Minimization or Maximization of Functions
- Finding roots can be tricky!

 $\diamond f(x) = 1 + |x|$ does not have any roots

 $\diamond 1 + x^2 \text{ and } \cosh(x)$ have roots, but not for real values of x

 $aggregation \sin(1/x)$ has ∞ number of roots – which one do you want?

- ◊ Make rough plot of function where (approximately) are the roots? which root do you want to find?
- We will consider the simplest algorithms

Bisection method

• Algorithm assumes

♦ f(x) changes sign at root (won't work for $f(x) = (x - 1)^2$) ♦ root is *bracketed* by x_0 and x_1 , i.e., $x_0 < \text{root} < x_1$ ♦ only one root in $[x_0, x_1]$

• Algorithm repeatedly bisects interval

$$\label{eq:compute } \begin{split} &\diamond \text{ compute } x_{\frac{1}{2}} = (x_0 + x_1)/2 \\ &\diamond \text{ compute product } f(x_0) \times f(x_{\frac{1}{2}}) \\ &\circ \text{ if product is positive replace } x_0 \leftarrow x_{\frac{1}{2}} \\ &\circ \text{ else replace } x_1 \leftarrow x_{\frac{1}{2}} \\ &\diamond \text{ repeat above steps until } |x_0 - x_1| \leq \epsilon \text{ (desired accuracy) or } \\ f(x_{\frac{1}{2}}) = 0 \end{split}$$

 \bullet The following program finds the root of $e^x \log(x) - x^2$

```
#include <cmath>
#include <iostream>
#include <iomanip>
using namespace std;
double f(double x) {
   return exp(x) * log(x) - x * x;
}
void print(int step, double x, double dx) {
    cout.setf(ios::right, ios::adjustfield);
    cout << " " << setw(4) << step << " ";
    cout.precision(15);
    cout.setf(ios::left, ios::adjustfield);
    cout.setf(ios::showpoint | ios::fixed);
   cout << setw(20) << x << " " << setw(20) << dx << '\n';
}
int main() {
   cout << " Bisection search for root of exp(x)*log(x) - x*x n"
        << " -----\n"
        << " Enter bracketing guesses x_0, x_1, and desired accuracy: ";
   double x0, x1, acc;
   cin >> x0 >> x1 >> acc;
```

```
dx\n"
cout << " Step
                          Х
                                                -----\n";
     << " ----
int step = 0;
double xHalf = (x0 + x1) / 2;
double dx = x1 - x0;
print(step, xHalf, dx);
double f0 = f(x0);
while (abs(dx) > abs(acc)) {
    double fHalf = f(xHalf);
    if (fHalf == 0) {
        dx = 0;
    } else {
        if (f0 * fHalf > 0) {
           x0 = xHalf;
           f0 = fHalf;
        } else {
           x1 = xHalf;
        }
        xHalf = (x0 + x1) / 2;
        dx = x1 - x0;
    }
    ++step;
    print(step, xHalf, dx);
}
```

}

Convergence rate

 \bullet Consider bisection after n steps

$$|dx_n| = |x_1 - x_0| \quad \text{after } n \text{ iterations} \\ = \frac{1}{2}|dx_{n-1}| = \frac{1}{2^2}|dx_{n-2}| = \dots = \frac{1}{2^n}|dx_0|$$

 \bullet Number of steps for accuracy ϵ

$$\frac{1}{2^n}|dx_0| \le \epsilon$$

. . .

solve for

$$n \ge \log_2 \left[\frac{|dx_0|}{\epsilon} \right] = \frac{\log_{10} \left[\frac{|dx_0|}{\epsilon} \right]}{0.3010 \dots}$$

Example: $|dx_0| = 0.1$ and $\epsilon = 10^{-6}$ requires $n \ge 17$

• General definition of convergence rate

$$|dx_n| \simeq C_F \, |dx_{n-1}|^{\alpha} \; ,$$

where order of convergence (= α) and convergence factor (= C_F)

Secant method

• For bisection $\alpha = 1$ – convergence is *linear* (rather slow)

Secant method

• Use *secant* (latin *secare* cut) line

$$s(x) = f(x_1) + \frac{f(x_1) - f(x_0)}{x_1 - x_0}(x - x_1)$$

• Next point determined by

$$s(x_{\text{new}}) = 0 \quad \Rightarrow \quad x_{\text{new}} = x_1 - (x_1 - x_0) \frac{f(x_1)}{f(x_1) - f(x_0)} \equiv x_1 + dx_{\text{new}} \;.$$

• Secant algorithm

♦ choose x_0, x_1 near root – need not bracket
♦ if $x(x_0) = f(x_1)$ algorithm fails – abort and retry
♦ otherwise replace $x_0 \leftarrow x_1$ and $x_1 \leftarrow x_{new}$ ♦ repeat above steps until $|x_0 - x_1| \le \epsilon$ (desired accuracy) or $f(x_{new}) = 0$

Secant method

• Convergence rate is the famous Golden Ratio

$$\alpha = \frac{1 + \sqrt{5}}{2} = 1.618033988\dots$$

• Convergence is *supelinear* – much faster than linear



• The following program finds the root of $e^x \log(x) - x^2$

```
#include <cmath>
#include <iostream>
#include <iomanip>
using namespace std;
double f(double x) {
   return exp(x) * log(x) - x * x;
}
void print(int step, double x, double dx) {
    cout.setf(ios::right, ios::adjustfield);
    cout << " " << setw(4) << step << " ";
    cout.precision(15);
    cout.setf(ios::left, ios::adjustfield);
    cout.setf(ios::showpoint | ios::fixed);
   cout << setw(20) << x << " " << setw(20) << dx << '\n';
}
int main() {
   cout << " Secant search for root of exp(x)*log(x) - x*x n"
        << " -----\n"
        << " Enter guesses x_0, x_1, and desired accuracy: ";
   double x0, x1, acc;
   cin >> x0 >> x1 >> acc;
```

}

```
dx\n"
cout << " Step
                X
    << " ----
                                             -----\n":
int step = 0;
double dx = x1 - x0;
print(step, x1, dx);
double f0 = f(x0);
while (abs(dx) > abs(acc)) {
   double f1 = f(x1);
   if (f0 == f1) {
       cerr << " Secant horizontal ... try again!\n";
       return 1;
   } else {
       dx *= -f1 / (f1 - f0);
       x0 = x1;
       f0 = f1;
       x1 += dx;
    }
   ++step;
   print(step, x1, dx);
}
```

Newton's tangent method

• Requires only one intial guess x_0 sufficiently close to root

Newton's tangent method

• Construct tangent (latin *tangere* touch) line

$$t(x) = f(x_0) + f'(x_0)(x - x_0)$$

 \bullet Next point is intersection with x axis

$$x_{\text{new}} = x_0 - \frac{f(x_0)}{f'(x_0)} \equiv x_0 + dx$$

- \bullet Convergence rate $\alpha=2$ is $\mathit{quadratic}$ very fast
- Requires derivative f'(x) either analytically or as finite difference

$$f'(x_0) \simeq \frac{f(x_1) - f(x_0)}{x_1 - x_0}$$

 \bullet The following program finds the root of $e^x \log(x) - x^2$

tangent.cpp

```
#include <cmath>
#include <iostream>
#include <iomanip>
using namespace std;
```

Newton's tangent method

```
double f(double x) {
   return exp(x) * log(x) - x * x;
}
double fPrime(double x) {
   return \exp(x) * (\log(x) + 1/x) - 2 * x;
}
void print(int step, double x, double dx) {
   cout.setf(ios::right, ios::adjustfield);
   cout << " " << setw(4) << step << " ";
   cout.precision(15);
   cout.setf(ios::left, ios::adjustfield);
   cout.setf(ios::showpoint | ios::fixed);
   cout << setw(20) << x << " " << setw(20) << dx << '\n';
}
int main() {
   cout << " Tangent search for root of exp(x)*log(x) - x*x n"
        << " -----\n"
        << " Enter guess x_0, and desired accuracy: ";
   double x, acc;
   cin >> x >> acc;
```

Computational Physics Library

}

```
dx\n"
cout << " Step
                        Х
    << " -----
                                  -----\n":
int step = 0;
double dx = 1;
print(step, x, dx);
while (abs(dx) > abs(acc)) {
   double slope = fPrime(x);
   if (slope == 0) {
       cerr << " Tangent horizontal ... try again!\n";</pre>
       return 1;
   } else {
       dx = -f(x) / slope;
       x += dx;
   }
   ++step;
   print(step, x, dx);
}
```

Computational Physics Library

• The root finding routines outlined above are available in the Computational Physics Library, see the header file findroot.hpp